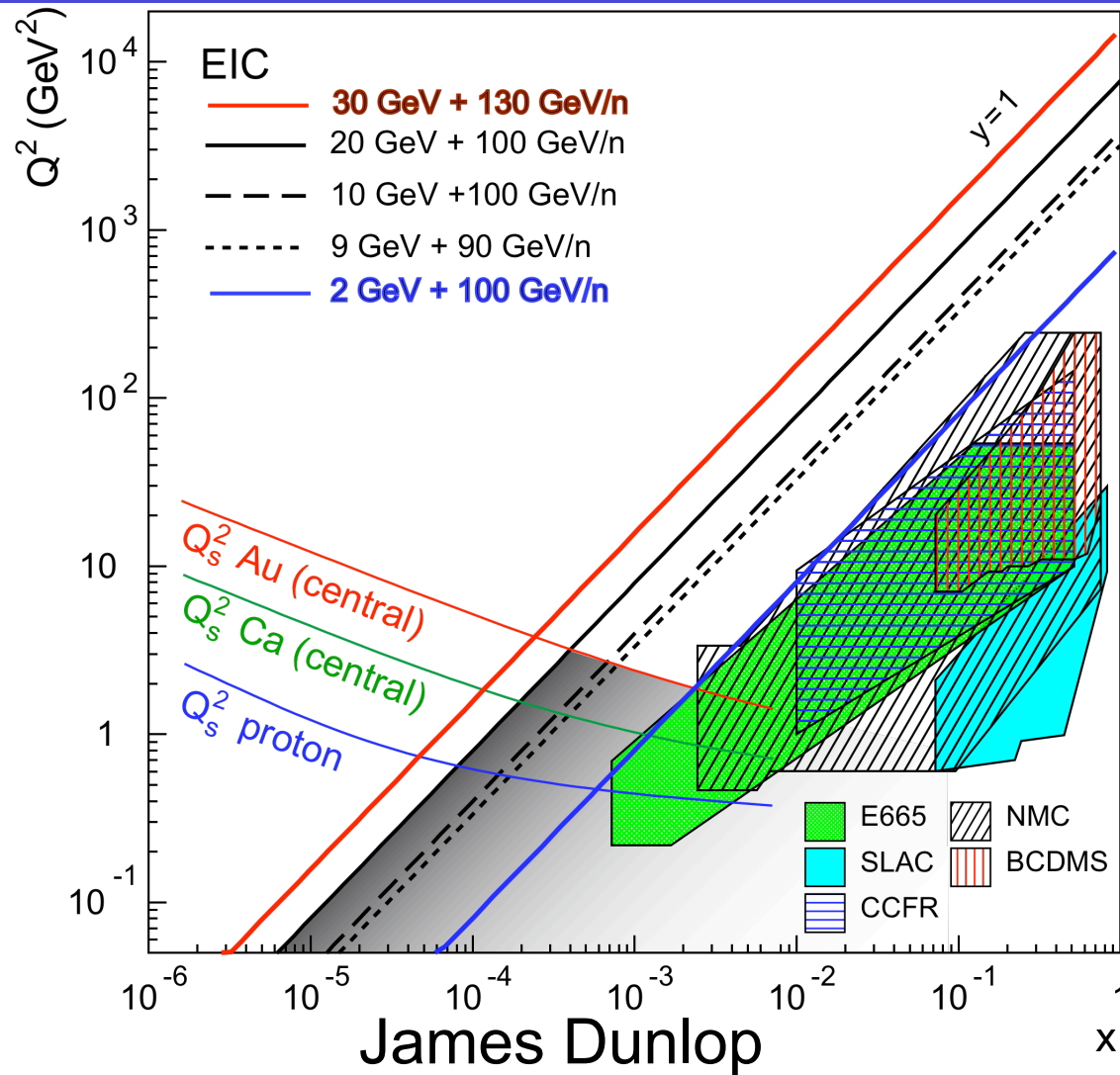


# e-A and the Low-x Structure of Matter



Brookhaven National Laboratory

# Theory of Strong Interactions: QCD

---

$$L_{QCD} = \bar{q}(i\gamma^\mu \partial_\mu - m)q - g(\bar{q}\gamma^\mu T_a q)A_\mu^a - \frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu}$$

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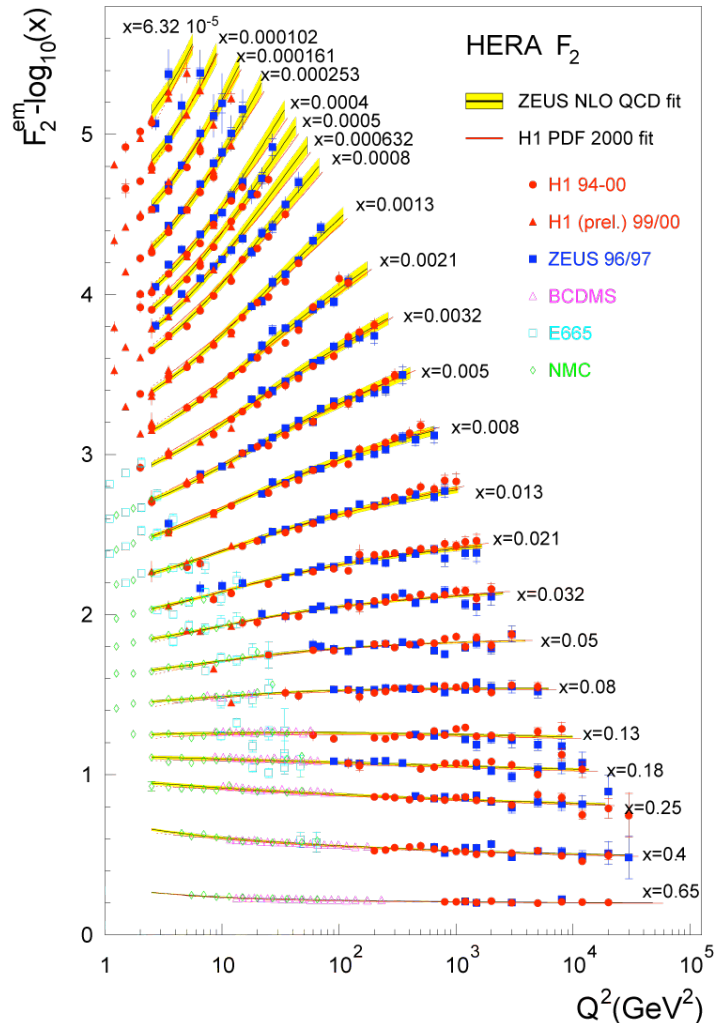
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⇒ QCD requires *fundamental* investigation via *experiment*

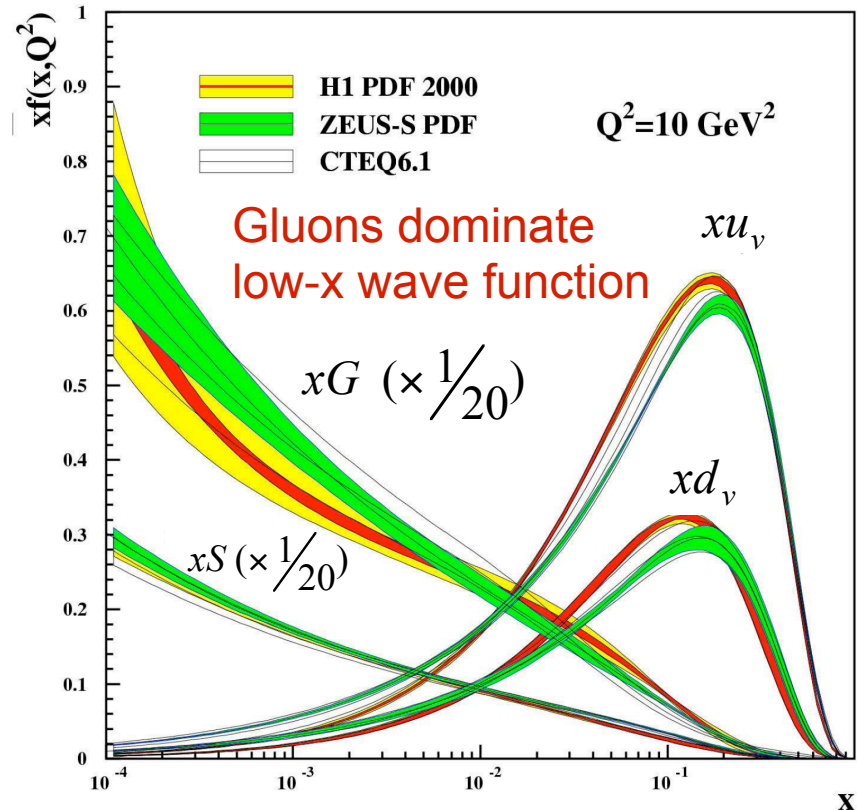


# What Do We Know About Glue in Matter?

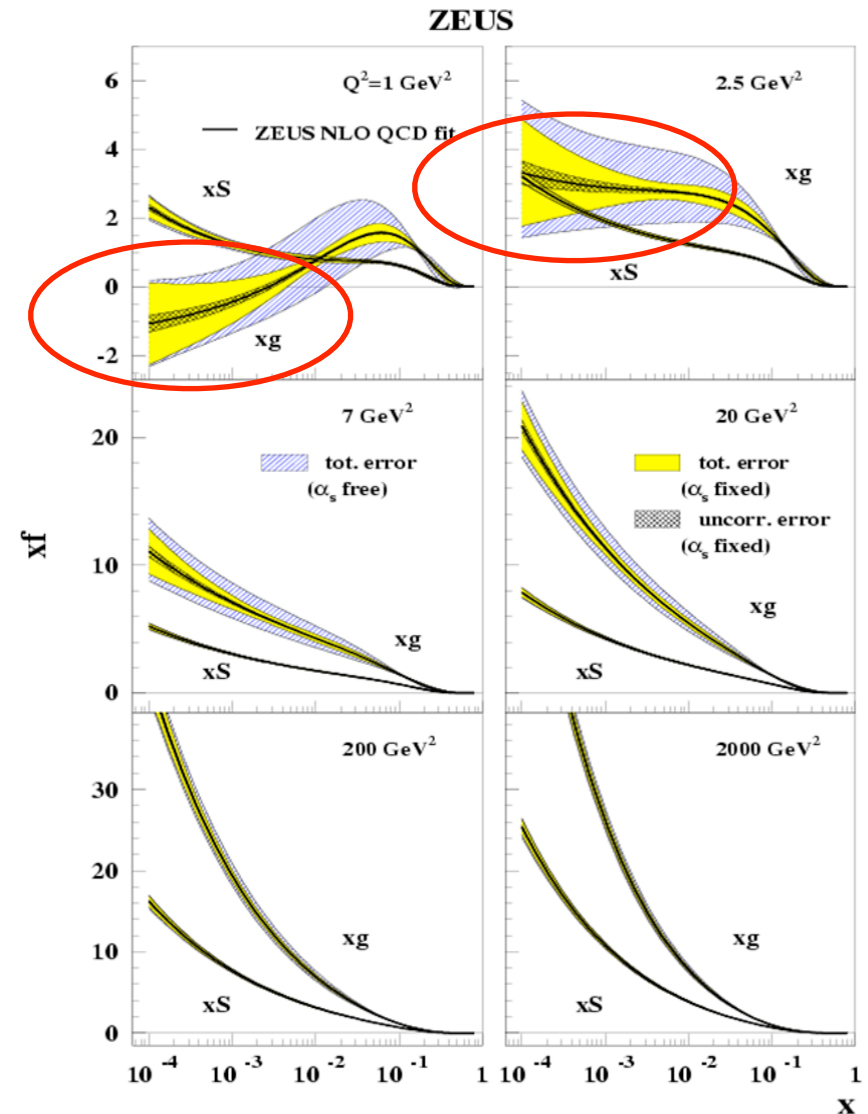
Deep Inelastic Scattering: 
$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha_{e.m.}^2}{xQ^4} \left[ \left( 1 - y + \frac{y^2}{2} \right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$



- Scaling violation:  $dF_2/d\ln Q^2$  and linear DGLAP Evolution  $\Rightarrow G(x, Q^2)$



# The Issue With Our Current Understanding



# The Issue With Our Current Understanding

## Established Model:

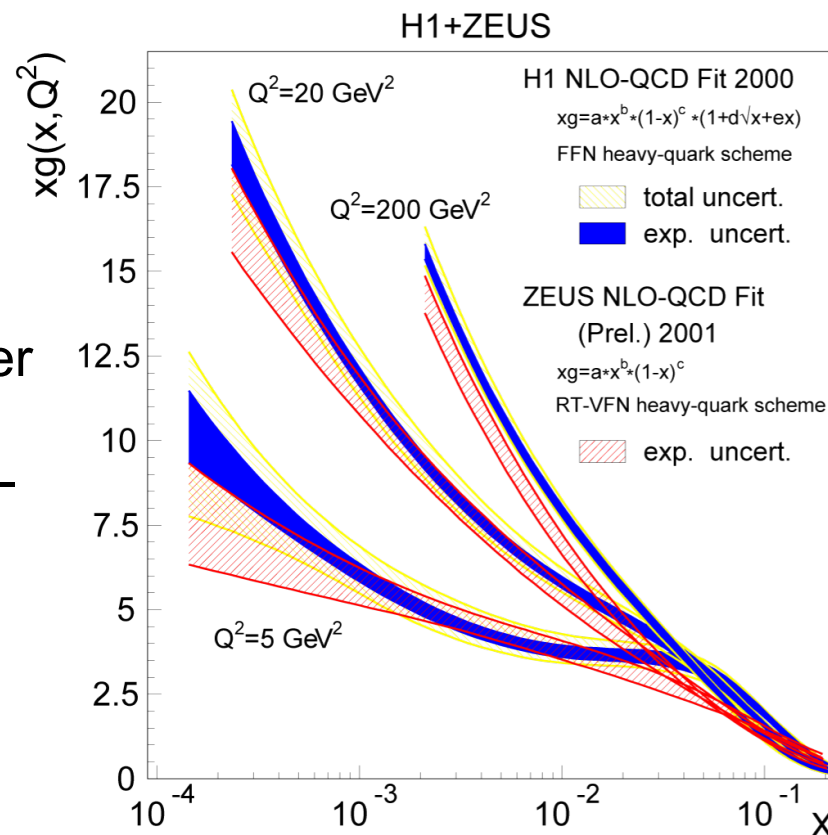
### Linear DGLAP evolution scheme

- Weird behavior of  $xG$  and  $F_L$  from HERA at small  $x$  and  $Q^2$ 
  - Could signal saturation, higher twist effects, need for more/better data?
- Unexpectedly large diffractive cross-section

more severe:

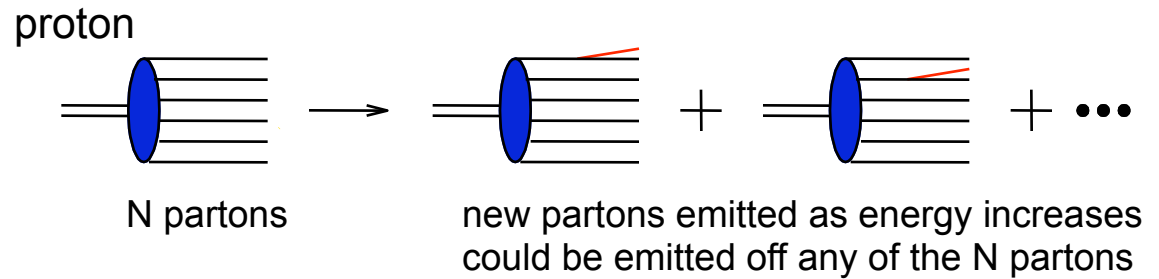
Linear Evolution has a built in high energy “catastrophe”

- $xG$  rapid rise for decreasing  $x$  and violation of (Froissart) unitary bound
- $\Rightarrow$  **must saturate**
  - What’s the underlying dynamics?

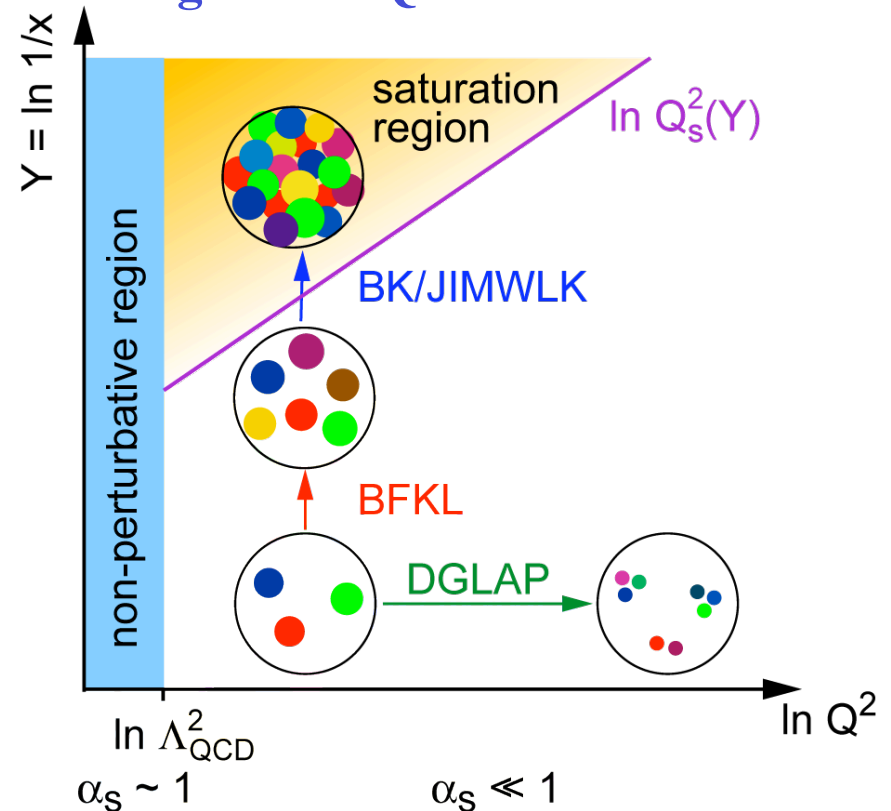


$\Rightarrow$  Need new approach

# Non-Linear QCD - Saturation



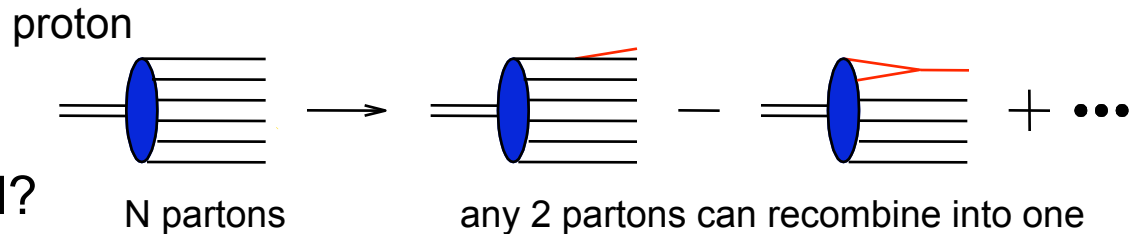
## Regimes of QCD Wave Function



# Non-Linear QCD - Saturation

- **BFKL Evolution in  $x$**

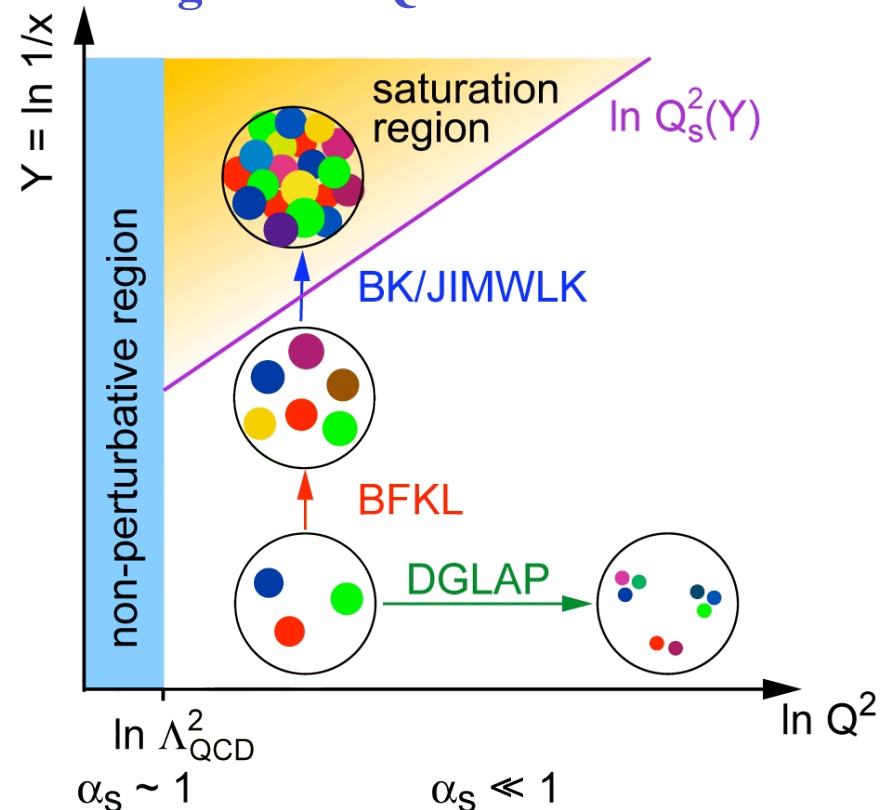
- linear
- explosion of color field?



- **New: BK/JIMWLK based models**

- introduce *non-linear effects*
- ⇒ saturation
- characterized by a scale  $Q_s(x, A)$
- arises naturally in the **Color Glass Condensate (CGC)** framework

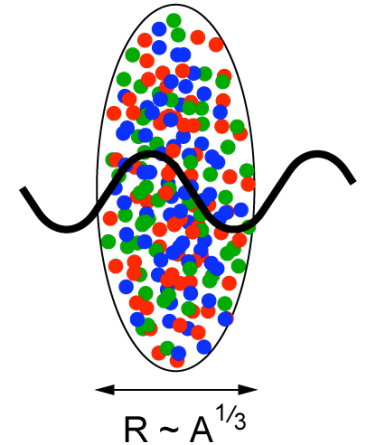
## Regimes of QCD Wave Function



# e+A: Studying Non-Linear Effects

## Scattering of electrons off nuclei:

- Probes interact over distances  $L \sim (2m_N x)^{-1}$
- For  $L > 2 R_A \sim A^{1/3}$  probe cannot distinguish between nucleons in front or back of nucleon
- Probe interacts *coherently* with all nucleons



$$Q_s^2 \sim \frac{\alpha_s x G(x, Q_s^2)}{\pi R_A^2} \quad \text{HERA: } xG \sim \frac{1}{x^{0.3}} \quad \text{A dependence: } xG_A \sim A$$

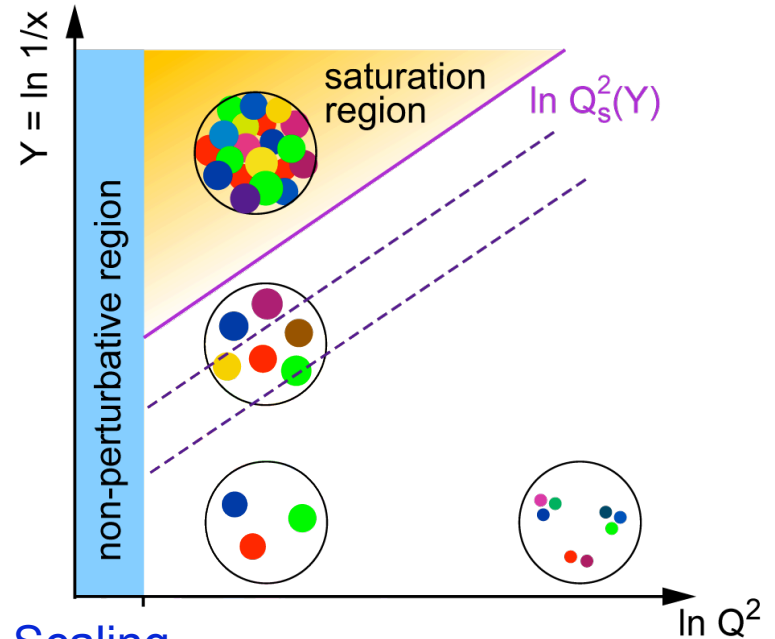
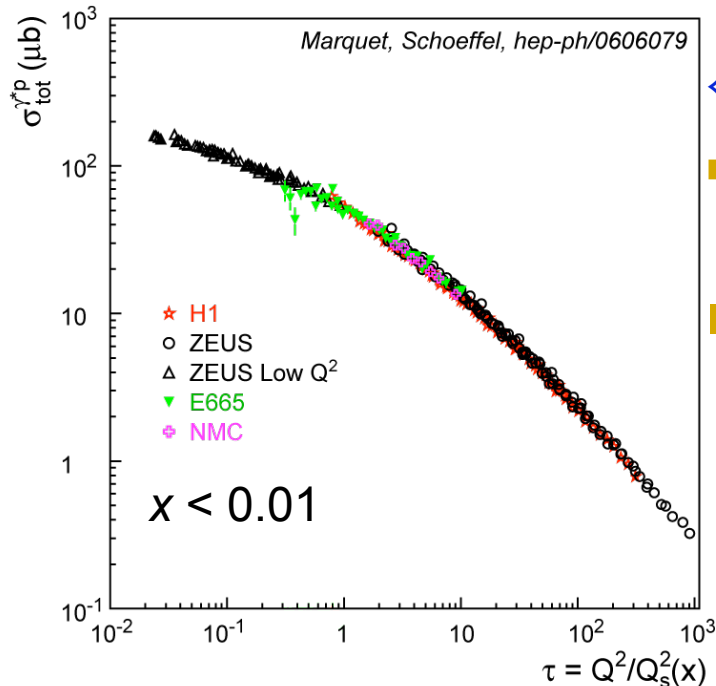
Nuclear “Oomph” Factor  
Pocket Formula:

$$(Q_s^A)^2 \approx c Q_0^2 \left( \frac{A}{x} \right)^{1/3}$$

Enhancement of  $Q_s$  with  $A \Rightarrow$  non-linear QCD regime reached at significantly lower energy in  $A$  than in proton

# Hints for Saturation at HERA & Geometric Scaling?

- Crucial *consequence* of non-linear evolution towards saturation:
- Physics *invariant* along trajectories parallel to saturation regime (lines of constant gluon occupancy)
- Scale with  $Q^2/Q_s^2(x)$  instead of  $x$  and  $Q^2$  separately



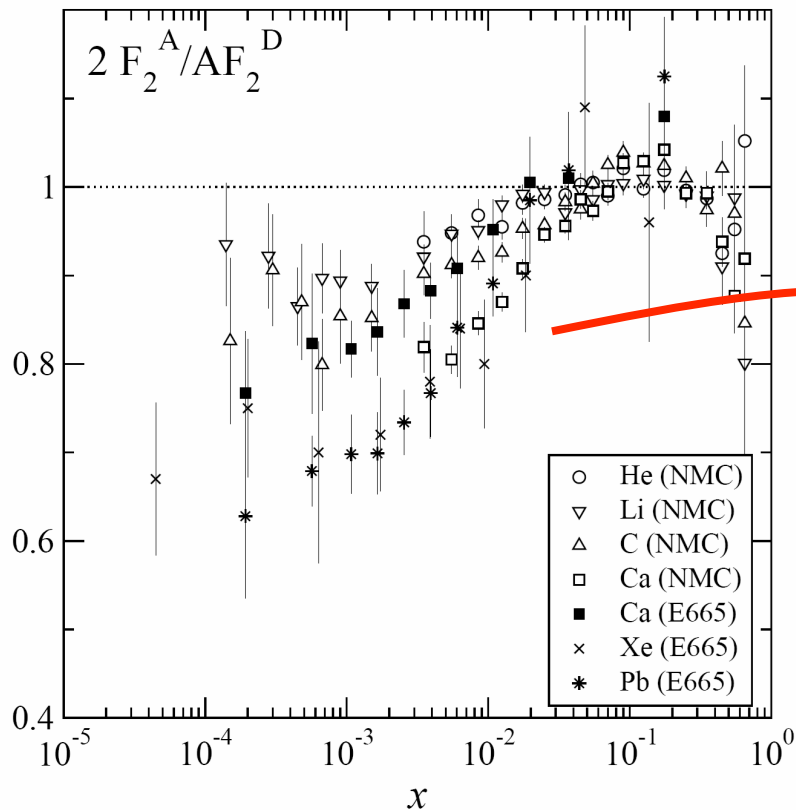
⇐ Geometric Scaling

- Consequence of saturation which manifests itself up to  $k_T > Q_s$
- Also seen in other final states (diffraction & VM production)

Scaling not proof but allows to set upper limit for saturation effects  
 $x < 10^{-2}$

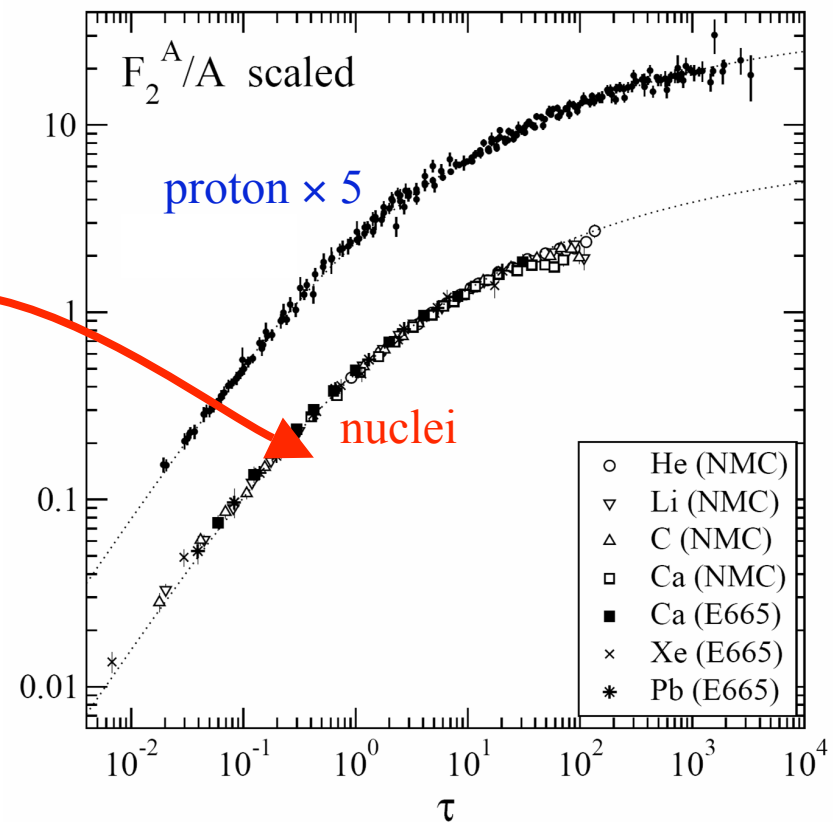
# Earlier Nuclear Experiments and Saturation

## Nuclear shadowing:



Freund et al., hep-ph/0210139

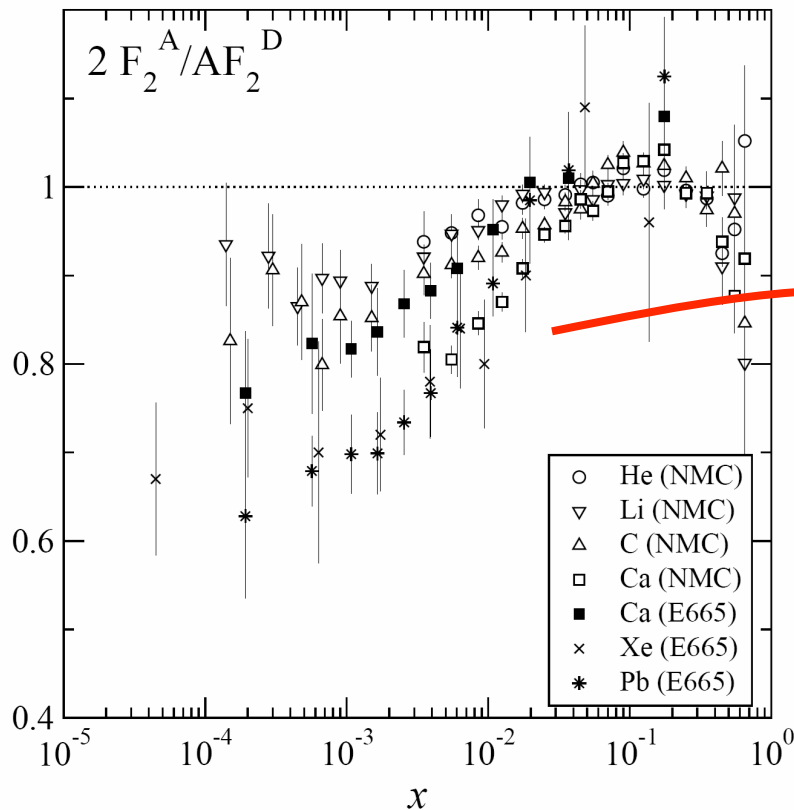
## Geometrical scaling





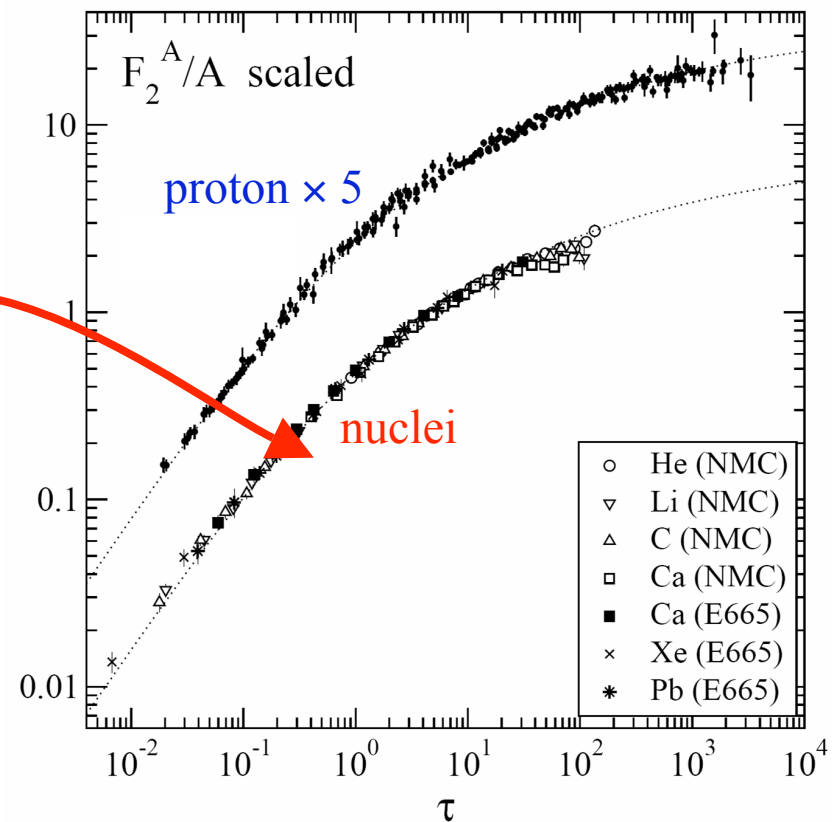
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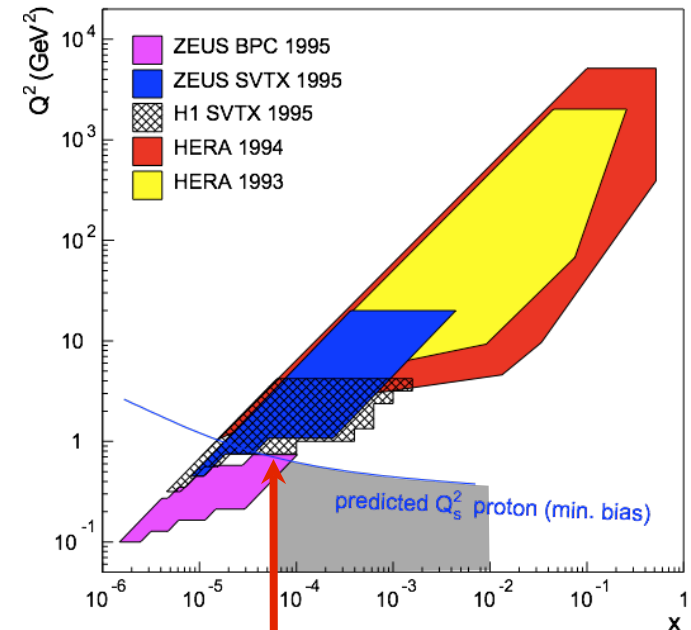
## Geometrical scaling



Geometrical scaling also found in nuclear experiments

# HERA & Saturation

- HERA (ep):
- Despite energy range far higher than EIC:
- $G_p(x, Q^2)$  through scaling violation known only outside (or in a very small region of) the saturation regime
- Same for  $G_p(x, Q^2)$  through  $F_L$
- HERA will provide a first direct measurement of  $G(x, Q^2)$  in the proton BUT
- Regime where non-linear QCD (saturation phenomena) matter ( $Q < Q_s$ ) out of reach!
- EIC: all relies on the Nuclear OOMPH (i.e. increasing  $Q_s$ )



# The Oomph Factor

- Nuclear Oomph Factor:  $(Q_s^A)^2 \approx c Q_0^2 \left(\frac{A}{x}\right)^{1/3}$

## Enhancement of $Q_s$ with $A$

⇒ non-linear QCD regime reached at significantly lower energy in  $e+A$  than in  $e+p$

$s_{Hera} \approx (330 \text{ GeV})^2$       Instead of extending  $x$ ,  $Q$  reach  
we increase  $Q_s$

$s_{EIC} \approx (63 \text{ GeV})^2$        $Q^2 \sim sx$ : EIC factor 27 behind  
(10+100 GeV)

$$\frac{s_{EIC}}{s_{Hera}} \approx \frac{1}{27}$$

$$Q_s^2(Hera) = Q_s^2(EIC) \rightarrow Q_0^2 x_{Hera}^{-1/3} = c Q_0^2 A^{1/3} x_{EIC}^{-1/3}$$

$$x_{EIC} = x_{Hera} \cdot c^3 A$$

$$c^3 A = 0.5^3 \cdot 197 \approx 25$$

# State-of-the-Art Oomph

- The e+A program lives and dies with the enhancement of  $Q_s^A$  over  $Q_s^p$

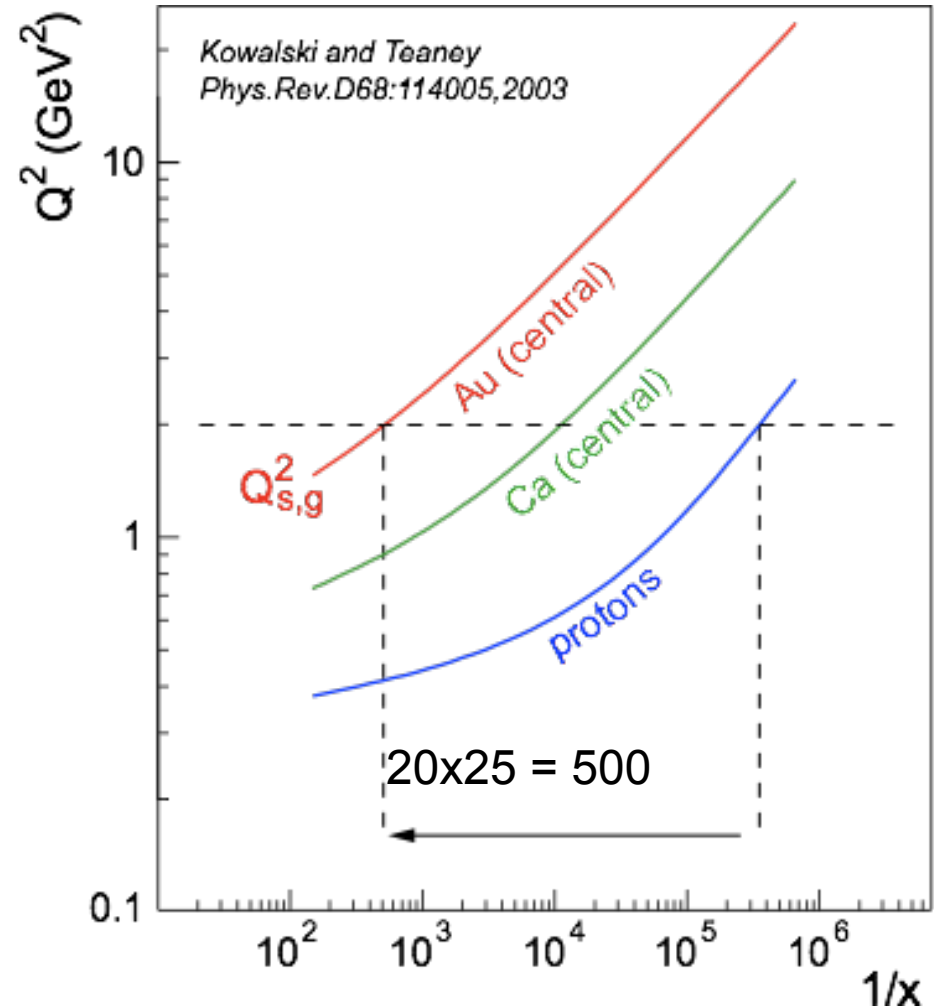
This factor is huge (500)

but

it's a model calculation!

- Assuming it's correct we "reach" further compared to HERA by  $500/27 = 18$
- (where we see no striking saturation effects)

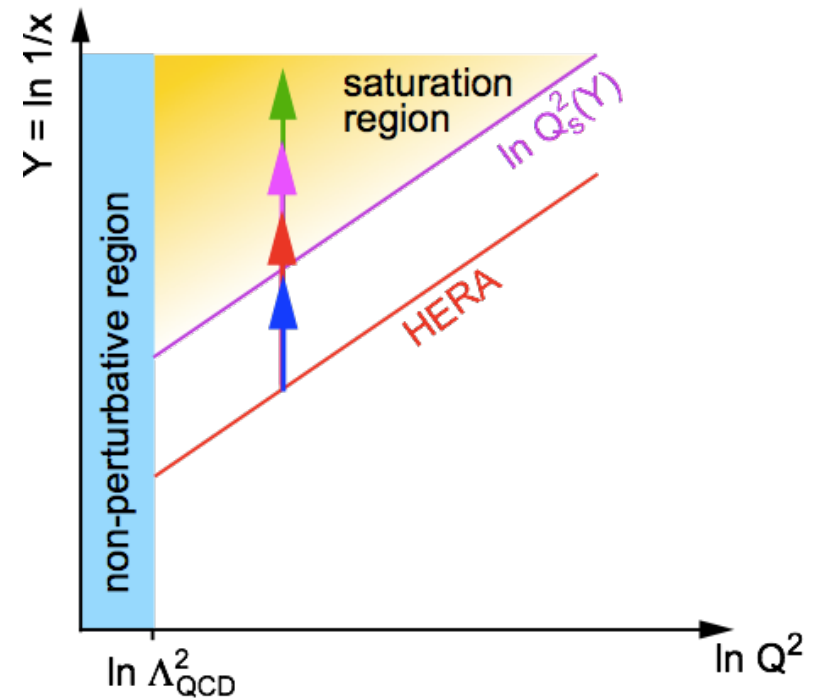
Here: protons for  $b=b_{\text{med}}$



# Reaching Saturation: Oomph versus HERA

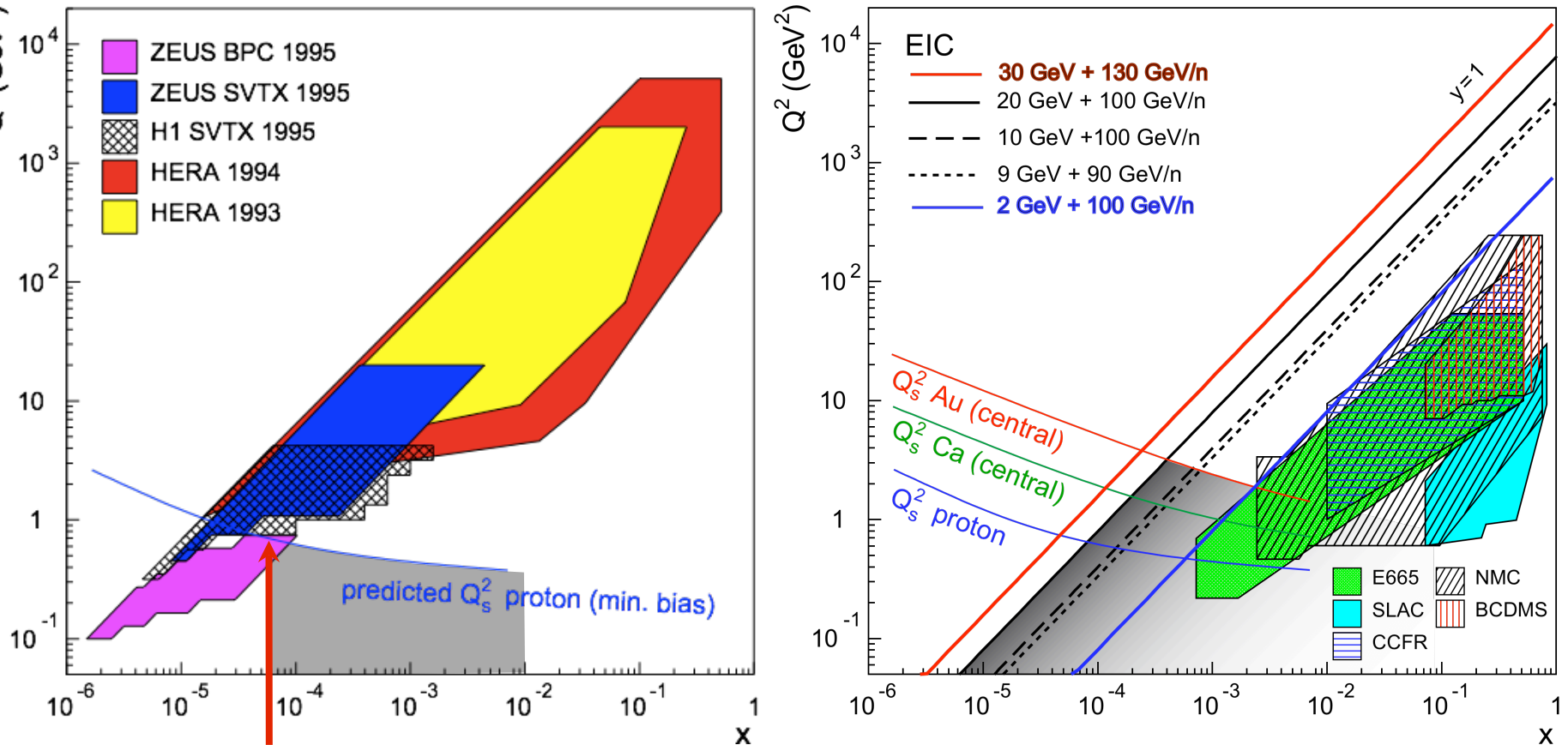
Beam Energy (GeV)	$\sqrt{s}$ (GeV)	SEIC/ SHERA	“virtual” x reach boost over HERA at $Q^2=\text{const}$
2+100	28	1/140	4
10+100	63	1/27	18
20+100	89	1/14	36
20+130	102	1/10	50
30+130	125	1/7	71

Numbers are rounded and approx. only



Note: We do not know (until we measured it) how far HERA was away from the saturation physics regime

# Reach compared with previous facilities



Staged option: begins to reach into the saturation regime for heavy nuclei

Experience with nuclei have shown that we need to reach deeply into a new regime for assurance that the new regime has been reached

And, we need a safety margin (models can and have been wrong before)

# Measurements: Understanding Glue in Matter

---

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---

Understanding the role of the glue in matter involves understanding its **key properties** which in turn define the required measurements:

- What is the momentum distribution of the gluons in matter?
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  - Strength of  $e+A$
- Do strong gluon fields effect the role of color neutral excitations (Pomerons)?
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# What is the Momentum Distribution of Gluons?

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## Gluon distribution $G(x, Q^2)$

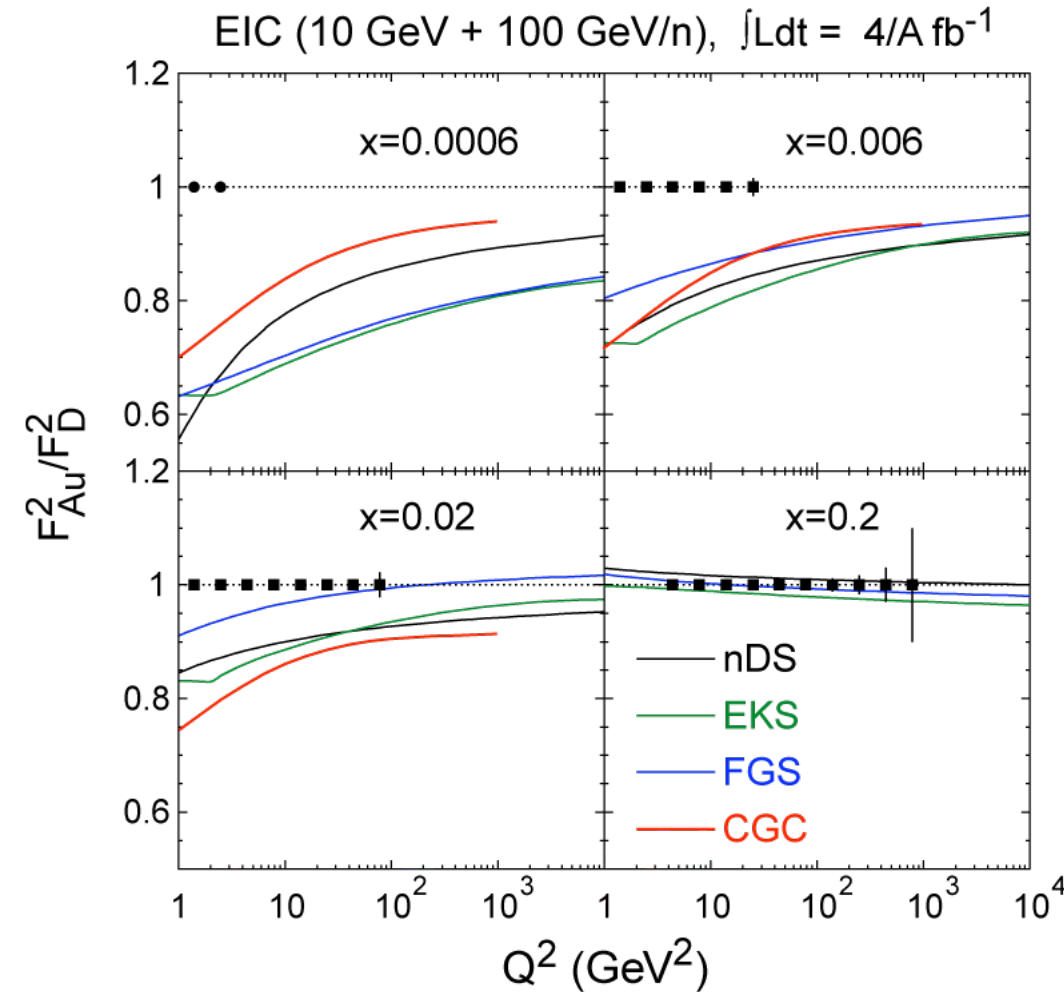
### – Shown here:

- Scaling violation in  $F_2$ :  $\delta F_2 / \delta \ln Q^2$
- $F_L \sim \alpha_s G(x, Q^2)$

### – Other Methods:

- 2+1 jet rates (needs jet algorithm and modeling of hadronization for inelastic hadron final states)
- inelastic vector meson production (e.g.  $J/\psi$ )
- diffractive vector meson production  $\sim [G(x, Q^2)]^2$ 
  - Active area of investigation
  - See M. Lamont's talk later today

# $F_2$ : Sea (Anti)Quarks Generated by Glue at Low $x$



$F_2$  will be one of the first measurements at EIC

nDS, EKS, FGS:

pQCD based models with different amounts of shadowing

Syst. studies of  $F_2(A, x, Q^2)$ :

$\Rightarrow G(x, Q^2)$  with precision

$\Rightarrow$  distinguish between models

$$\frac{d^2 \sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} \left[ \left( 1 - y + \frac{y^2}{2} \right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

# $F_L$ : measure glue directly

$$\frac{d^2\sigma^{ep \rightarrow eX}}{dx dQ^2} = \frac{4\pi\alpha^2}{xQ^4} \left[ \left( 1 - y + \frac{y^2}{2} \right) F_2(x, Q^2) - \frac{y^2}{2} F_L(x, Q^2) \right]$$

$$F_L \sim \alpha_s G(x, Q^2)$$

requires  $\sqrt{s}$  scan,  $Q^2/xs = y$

• Assume:

•  $L = 3.8 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

•  $T = 10 \text{ weeks}$

• duty cycle: 50%

•  $L \sim 1/A$  (approx)

•  $\int L dt = 11 \text{ fb}^{-1}$

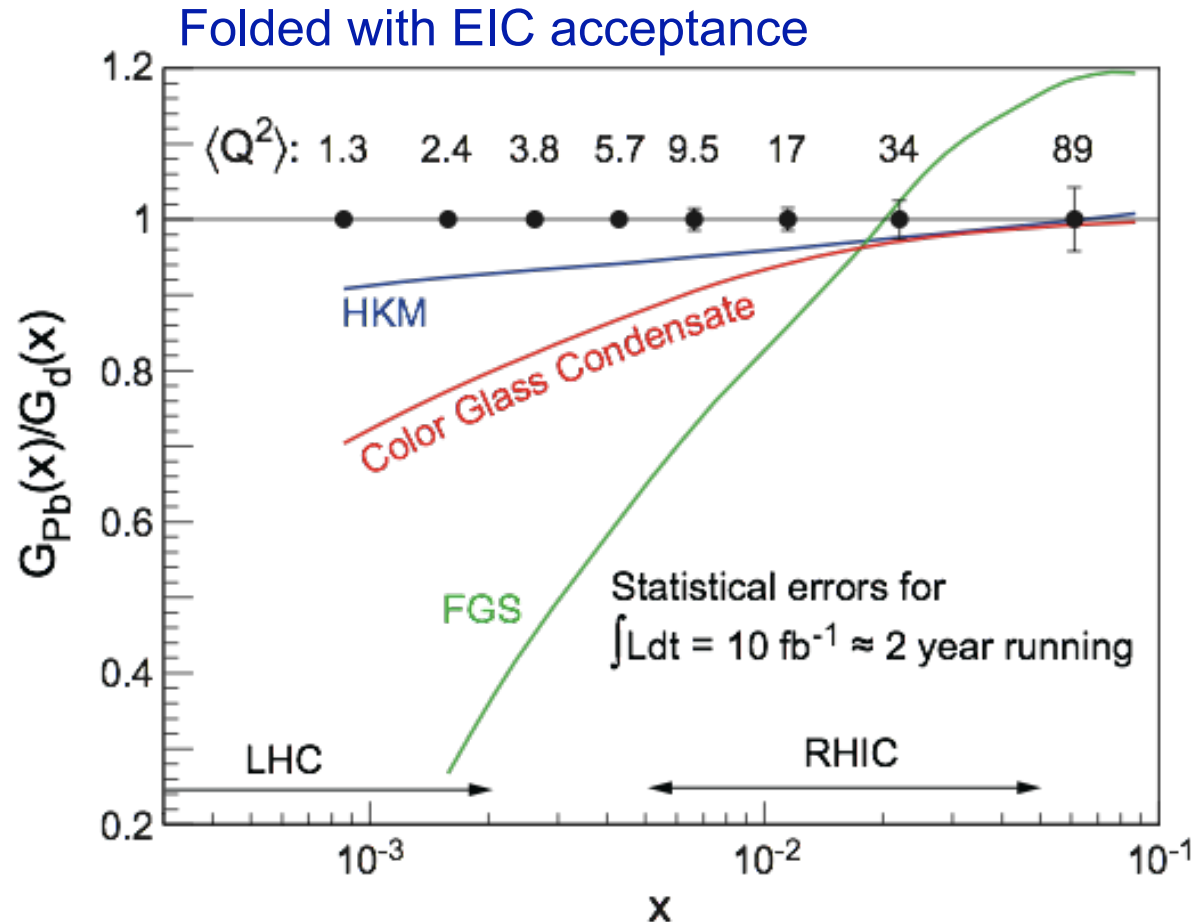
Plot contains:

$\int L dt = 4/A \text{ fb}^{-1} \text{ (10+100) GeV}$

$= 4/A \text{ fb}^{-1} \text{ (10+50) GeV}$

$= 2/A \text{ fb}^{-1} \text{ (5+50) GeV}$

statistical error only



# Existing $F_L$ measurements: Hera

## •Runs in 2007:

- High  $s$  -  $E_p=920$  GeV ( $H1 = 21.9 \text{ pb}^{-1}$ )
- Low  $s$  -  $E_p=460$  GeV ( $H1 = 12.4 \text{ pb}^{-1}$ )
- Medium  $s$  -  $E_p=575$  GeV ( $H1 = 6.2 \text{ pb}^{-1}$ )

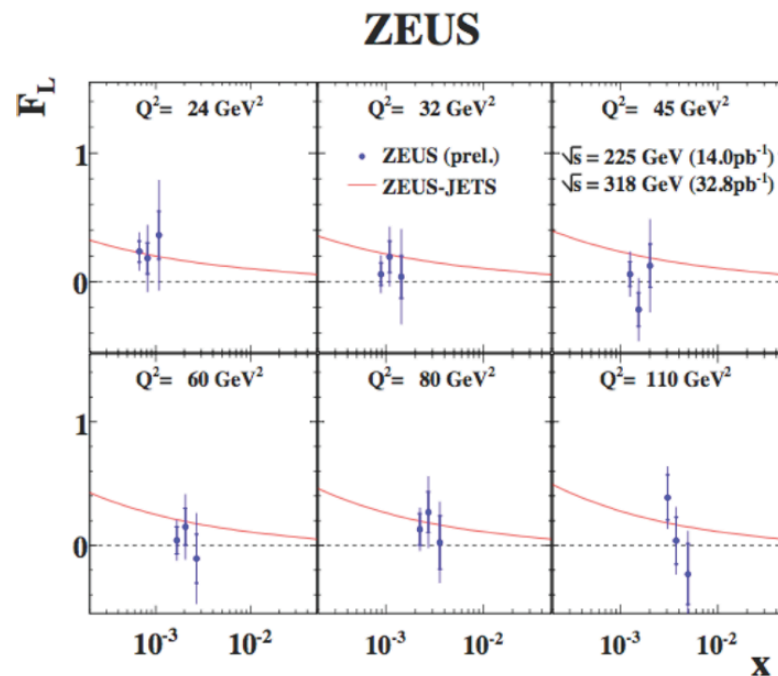
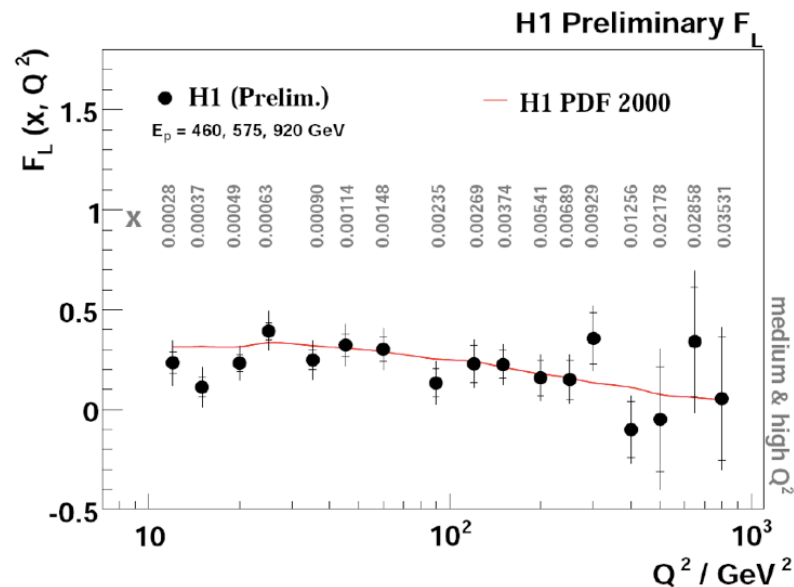
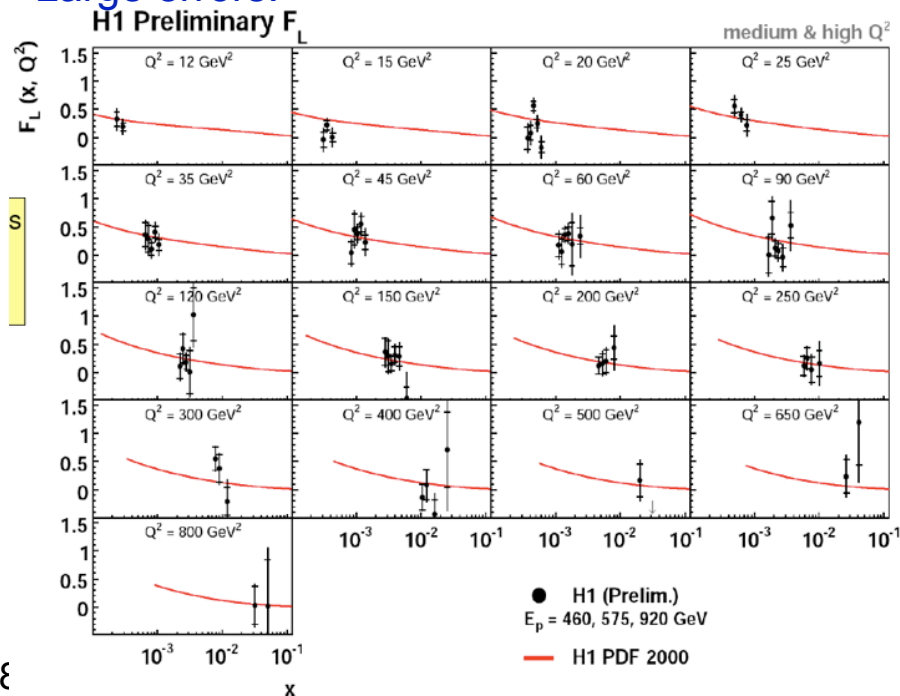
## •Sensitivity to $F_L$ requires high $y$

## •Challenge - high $y$ means low electron energy

## •HERA: $F_L(H1) > F_L(\text{Zeus})$ ?

limited to large  $Q$  ( $Q > Q_s$ )  $x = 10^{-4} \dots 10^{-1}$

## •Large errors!





# $F_L$ and Syst. Errors

- W/o at least a rough detector design and lots of simulations it is hard to estimate sys. uncertainties

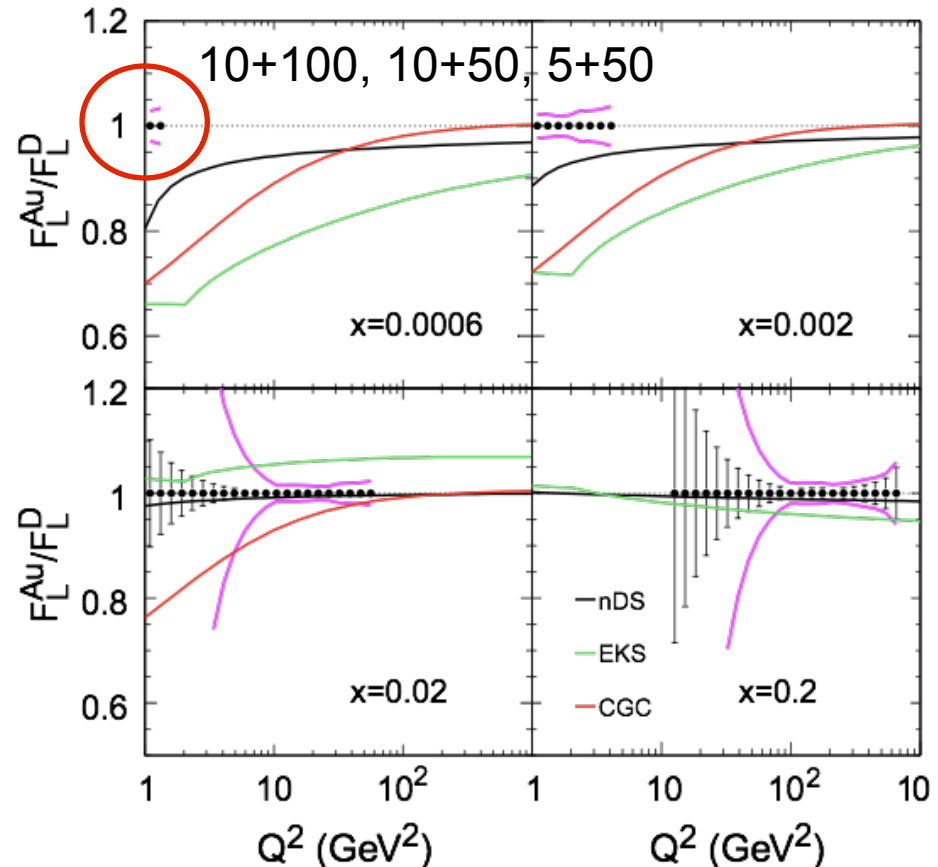
Simple estimate J. Dunlop/A. Bruell:  
1% energy-to-energy normalization (only)  
following discussions at MIT EIC Mtg.

- How realistic are the assumptions?
- Compare to current HERA studies ?

Conclusion from [this study](#):

**Dominated by sys. Uncertainties**

Luminosity not the limit, but need  
more detailed studies (w detector)



Need to maximize  $y$  range, maximize range of  $s$  scanned

e.g.  $x=0.005$ ,  $Q^2 = 2$  GeV<sup>2</sup>:  $y$  from 0.5 (2+100) to 0.03 (30+130)

# The Gluon Space-Time Distribution

---

- What we know is mostly the momentum distribution of glue
  - How is the glue distributed spatially in nuclei?
  - Gluon density profile: small clumps or uniform ?
- Various techniques & methods:
  - Exclusive final states (e.g. vector meson production  $\rho$ ,  $J/\psi$ , DVCS)
    - color transparency  $\Leftrightarrow$  color opacity
  - Deep Virtual Compton Scattering (DVCS)
    - Integrated DVCS cross-section:  $\sigma_{\text{DVCS}} \sim A^{4/3}$
  - Measurement of structure functions for various mass numbers  $A$  (shadowing, EMC effect) and its impact parameter dependence
- Promising direction: fundamentally new approach in nuclei from which much can be learned even at the lower energies

# Hadronization and Energy Loss

## nDIS:

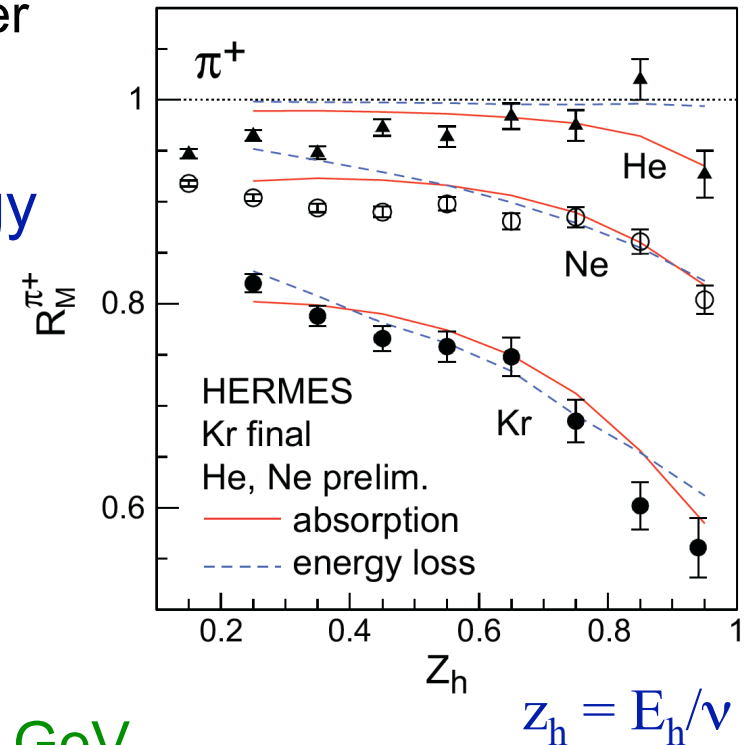
- Suppression of high- $p_T$  hadrons analogous but *weaker* than at RHIC
- Clean measurement in ‘cold’ nuclear matter

## Fundamental question:

What is the mechanism for QCD energy loss in matter?

When do colored partons get neutralized?

Parton energy loss vs.  
(pre)hadron absorption



Energy transfer in lab rest frame

EIC:  $10 < \nu < 1600$  GeV    HERMES: 2-25 GeV

EIC: can measure *heavy flavor* energy loss

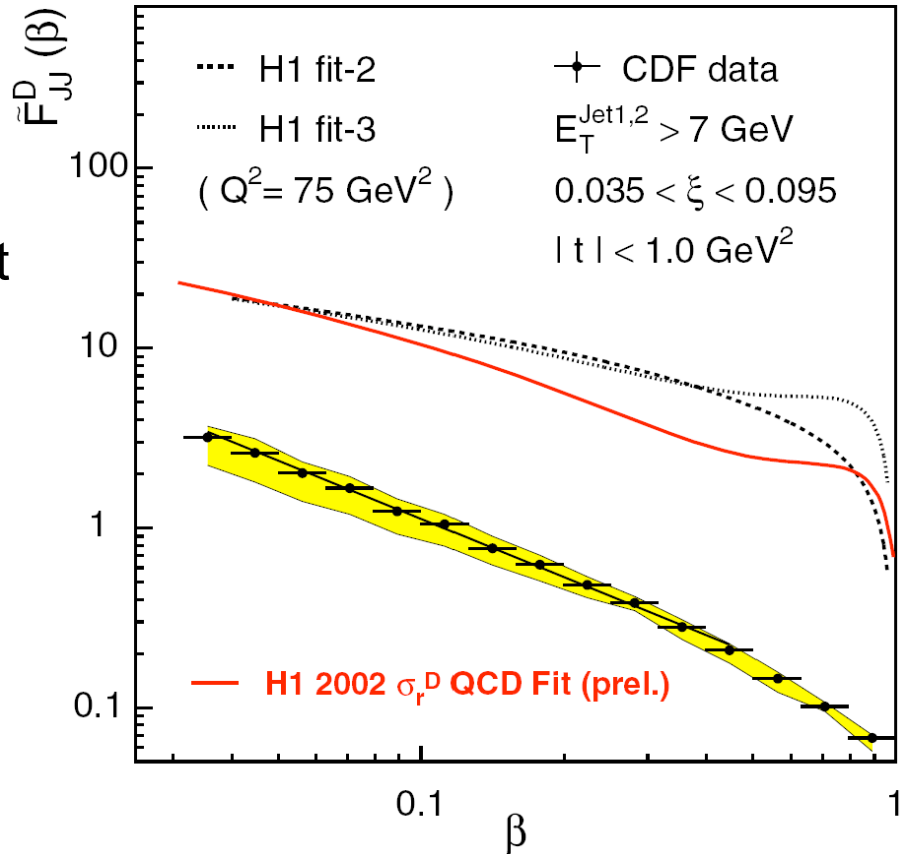
Mass effects not understood at RHIC, control time scales

x range not required to be small, can start at 2+100

# Connection to $p+A$ Physics

- $e+A$  and  $p+A$  provide excellent information on properties of gluons in the nuclear wave functions
- Both are **complementary** and offer the opportunity to perform stringent checks of **factorization/universality**

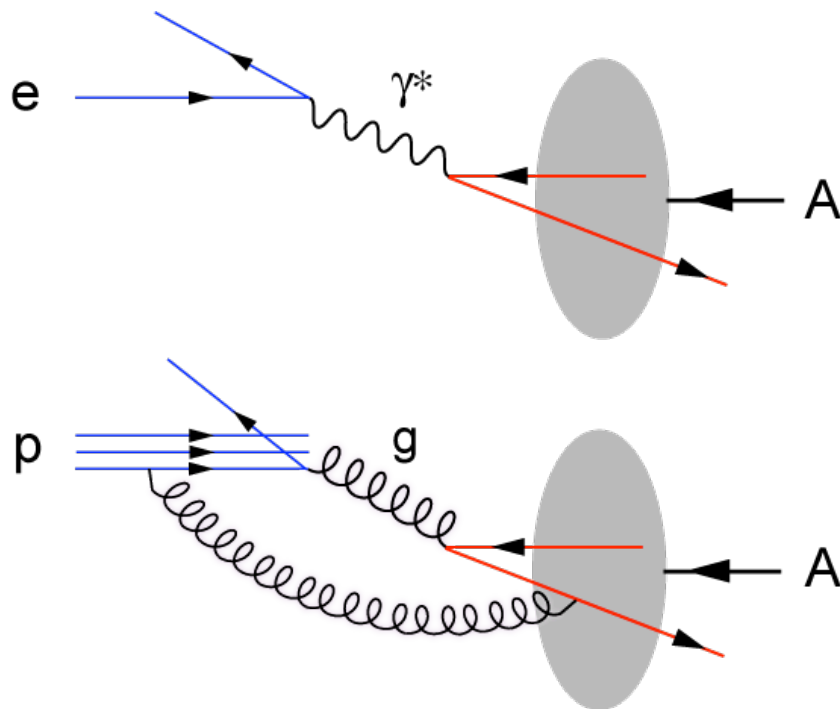
F. Schilling, hex-ex/0209001



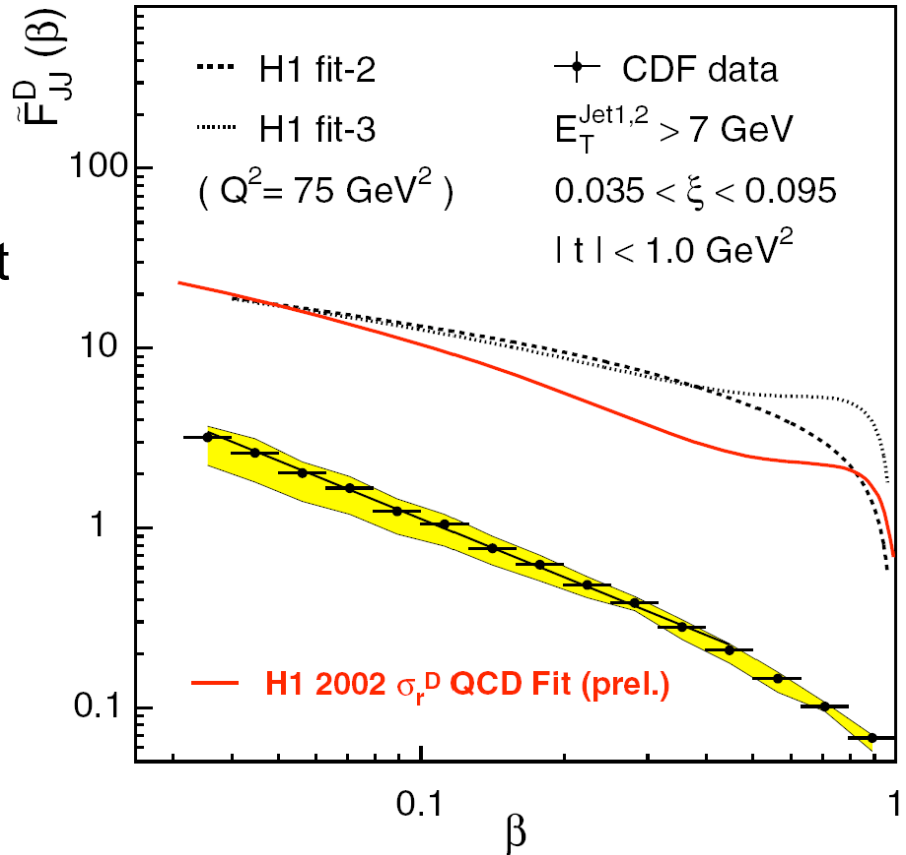
**Breakdown** of factorization ( $e+p$  HERA versus  $p+p$  Tevatron) seen for diffractive final states.

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- Both are **complementary** and offer the opportunity to perform stringent checks of **factorization/universality**



F. Schilling, hex-ex/0209001



**Breakdown** of factorization ( $e+p$  HERA versus  $p+p$  Tevatron) seen for diffractive final states.

# Connection to RHIC & LHC Physics

## Matter at RHIC:

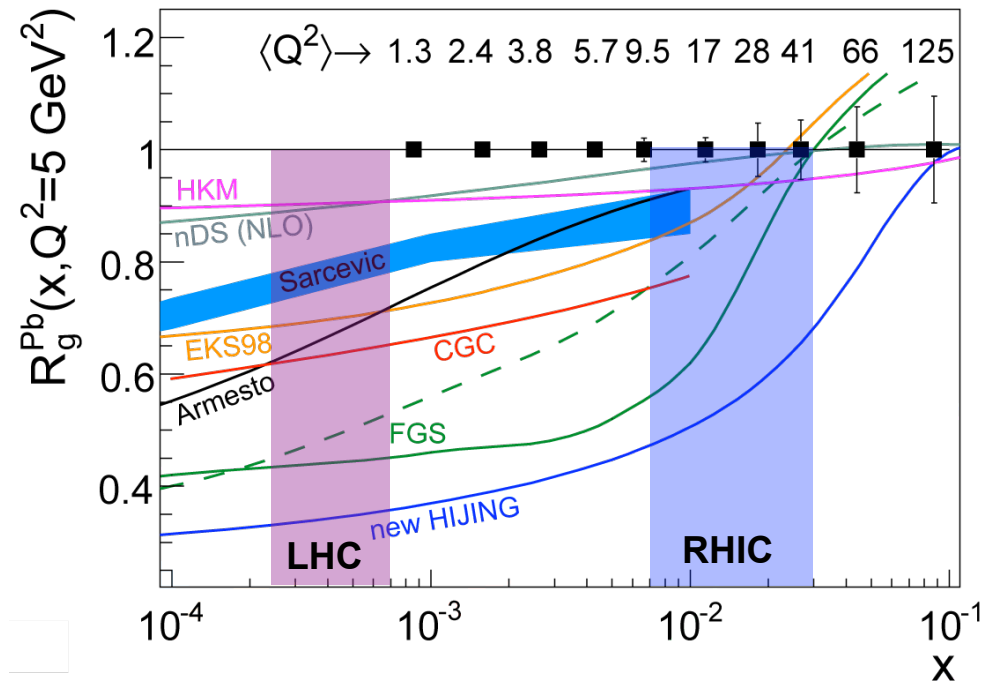
- thermalizes fast ( $\tau_0 \sim 0.6$  fm/c)
- We don't know why and how?
- Initial conditions?  $\Rightarrow G(x, Q^2)$

## Role of saturation ?

- RHIC  $\rightarrow$  forward region
- LHC  $\rightarrow$  midrapidity
  - bulk (low- $p_T$  matter) & semi-hard jets

## Jet Quenching:

- Need Reference: E-loss in cold matter
- No HERMES data for
  - charm energy loss
  - in LHC energy range

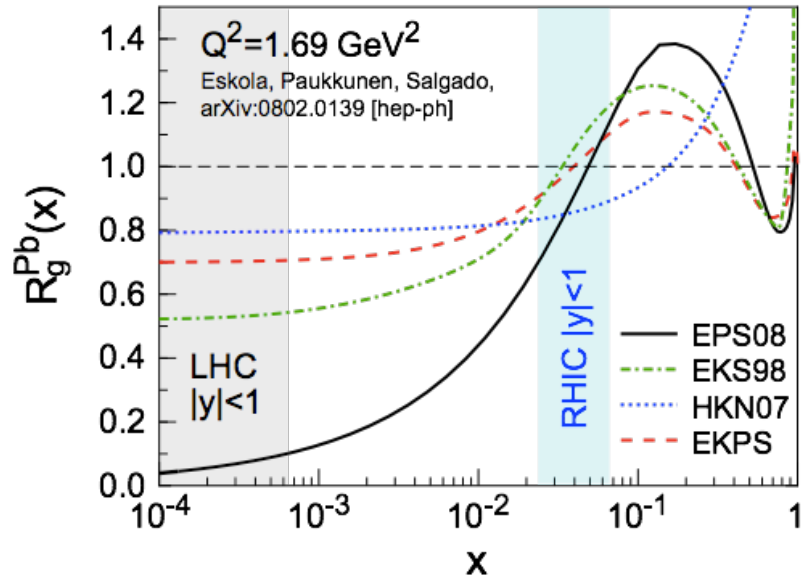


**EIC provides new essential input:**

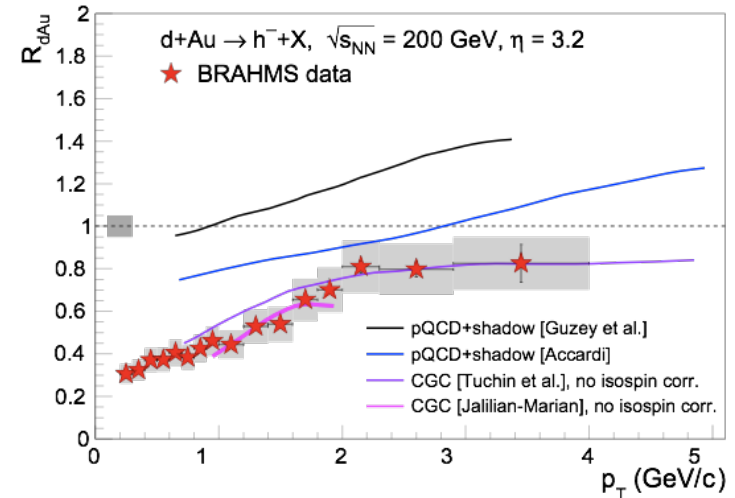
- Precise handle on  $x, Q^2$
- Means to study exclusive effects

# Connections with RHIC and LHC

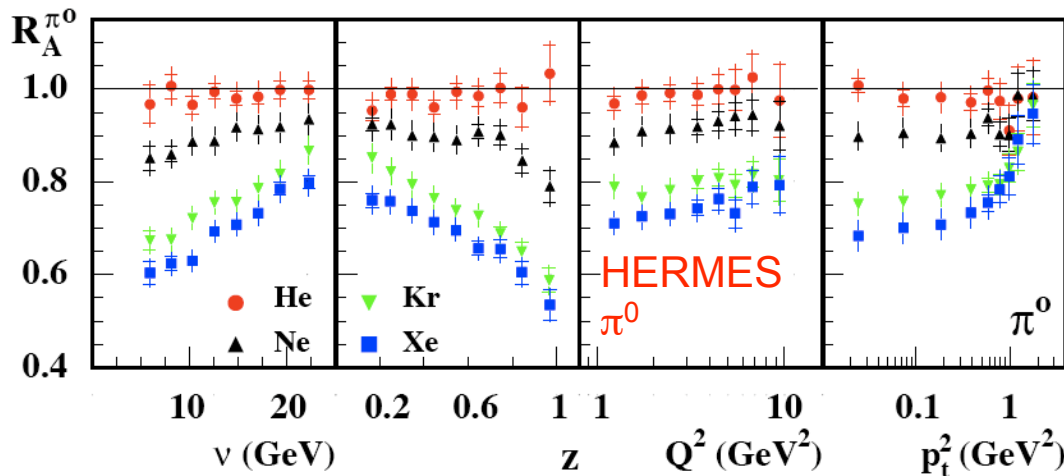
## Shadowing|Antishadowing|EMC



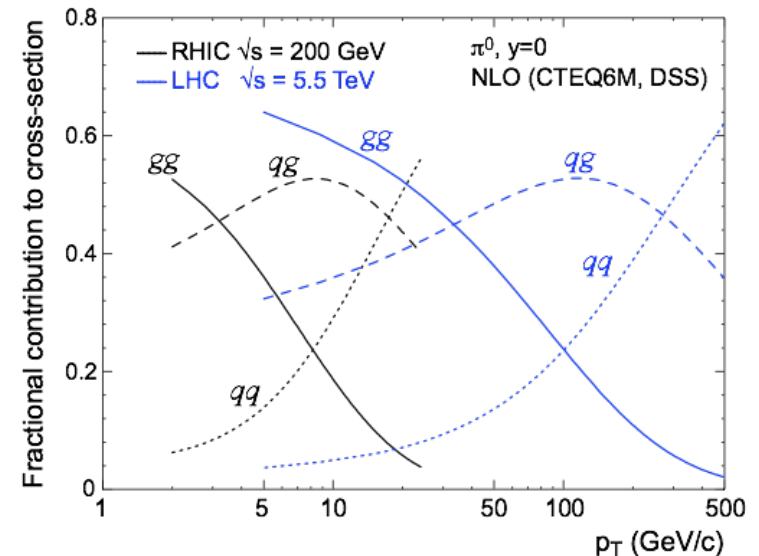
## Saturation (initial state) effects (RHIC fwd, LHC mid-rapidity)



## Hadron attenuation|Eloss in cold matter



## Particle Production



# Summary

EIC provides a chance to dive deeply into a fundamentally new regime of one of the four basic forces, QCD

## Issues:

Need to broaden and deepen measurements

Diffraction

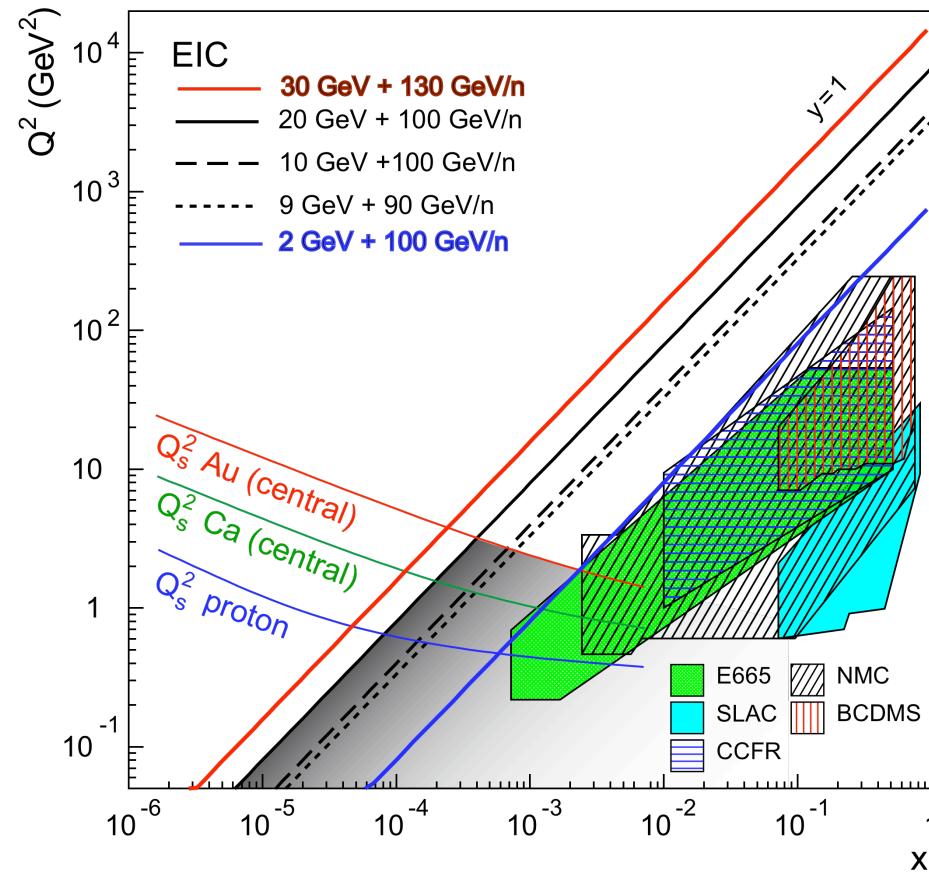
Jet-medium interactions

...

Need to develop connections

To RHIC/LHC

To larger scientific community



What is the smoking gun for crossing the saturation scale?